

# X-ray Emission Spectroscopy in Magnetic 3d-Transition Metals

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# X-rayEmission SpectroscopyinMagnetic 3d-TransitionMetals

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### **Introduction:**

The application of high pressure affects the band structure and magnetic interactions in solids by modifying nearest -neighbor distances and interatomic potentials. While all materials experience electronic changes with increasing pressure, spin polarized, strongly electron correlated materials are expected to undergo the most dramatic transformations. In such materials, (d and f-electron metals and compounds), applied pressure reduces the strength of on d-site correlations, leading to increase delectron delocalization and, eventually, to loss of its magnetism. In this ongoing project, we study the electronic and magnetic properties of Group VIII, d (Fe, Co and Ni) magnetic transition metals and their compounds at high pressures.

The high-pressure properties of magnetic 3d-transition metals and compounds have been studied extensively over the years, because of iron being a major constituent of the Earth's core and its relevance to the planetary modeling to understand the chemical com position, internal structure, and geomagnetism. However, the fundamental scientific interest in the high -pressure properties of magnetic 3d -electron systems extends well beyond the geophysical applications to include the electron correlation -driven physics. The role of magnetic interactions in the stabilization of the "non -standard" ambient pressure structures of Fe, Co and Ni is still incompletely understood. Theoretical studies have predicted (and high orrelationsbetweentheelectronicstructureandphase pressureexperiments are beginning to show) strong c stability in these materials. The phase diagrams of magnetic 3 d systems reflect a delicate balance between the diagram of tspin interactions and structural configuration. At ambient conditions, the crystal structures and  $\varepsilon$ -Co(hcp) phases depart from the standard sequence (hcp  $\rightarrow$ bcc $\rightarrow$ hcp $\rightarrow$ fcc), as observed in all other non-magnetic transition metals within creasing the d -bandoccupancy, and are different from those of their 4d- and 5d-counter parts. This anomal ous behavior has been interpreted in terms of the spin -polarizedd band altering the d -band occupancy [1]. At high pressures, however, the d -valence band is expected to broaden resulting in a suppression or even a complete loss of magnetism. Experimentally , ferromagnetic α(bcc)-Fehasbeenconfirmedtotransformtonon -magnetic ε–Fe (hcp)at10GPa[2,3].Recently, wehave also observed a similar transition in Cofrom ferromagnetic $\alpha$ (hcp)-Cotolikelynonmagnetic  $\beta$ (fcc)-Coat 105 GPa[4]. A similar structural phase transit ion is expected in Ni, probably in the second -order fcc -fcc transition. However, there has been no directly measured change in magnetism associated with the structural phase transition in Co, nor has yet been confirmed such an iso-structuralphasetransit ioninNi.

Similar electronic transitions have been proposed in these 3d -transition metal oxides (FeO, CoO and NiO) from high spin (magnetic) to low spin (nonmagnetic) states [5]. In each of these systems, the

magnetic transition is accompanied by a firs t-order structural transition involving large volume collapse (10% in FeO, for example). So far, there have been no electronic measurements under pressure confirming these significant theoretical predictions, although the predicted pressures for the volume ecollapse transitions are within the experimental pressure range (80 -200 GPa).

### MethodsandMaterials:

In our experiments, we use synchrotron based X-ray emission spectroscopy (XES) of the K  $\beta$  transition to determine the spin state of 3d  $\beta$  -transition metals and transition metaloxides under pressure. In K $\beta$  XES, 1s core -electrons are excited by high  $\beta$  -energy polychromatic x  $\beta$ -rays from a highly collimated synchrotron source (13ID  $\beta$ -D, GSE  $\beta$ -CARS). There sulting K  $\beta$  emission is analyzed using a high  $\beta$ -resolution (1eV) analyzer/detector system to determine the polarization of the valence band electrons. For the 3d transition elements, the K  $\beta$  line (3p  $\beta$ -1s transition) contains information about the polarization of the valence band via the 3p  $\beta$ -3d exchange interaction. The separation,  $\beta$ -3d exchange interaction. The separation,  $\beta$ -3d exchange interaction about the polarization of the valence band via the 3p  $\beta$ -3d exchange interaction. The separation,  $\beta$ -3d exchange interaction about the polarization of the valence band via the 3p  $\beta$ -3d exchange interaction. The separation  $\beta$ -3d exchange interaction about the polarization of the valence band electrons in  $\beta$ -5d exchange interaction. The separation  $\beta$ -5d exchange interaction about the polarization of the valence band electrons in  $\beta$ -5d exchange interaction. The separation  $\beta$ -5d exchange interaction about the polarization of the valence band electrons in  $\beta$ -5d exchange interaction. The separation  $\beta$ -6d exchange interaction about the polarization of the valence band electrons in  $\beta$ -5d exchange interaction.

Traditionally, the limiting factor in the DAC  $\,$ -XES experiments in 3d elements is that the K emission photons (7 -8 keV) are strongly attenuated by the diamond anvils. In addition, the small band splitting requires the use of high  $\,$ -intensity x -rays from collimated 3  $\,$ <sup>rd</sup> generation synchrotron sources, unavailable until recently for use in the DAC. To avoid the absorption of XES signal by the diamond anvils, our XES measurements have been conducted in a radial configuration, using hardened Be gaskets [7]. This allows simultaneous access along the axial direction, for use for in  $\,$ -situlaser -heating, on -line ruby pressuremeasurements, or  $\,$ -ray diffraction measurements.

### **Results:**

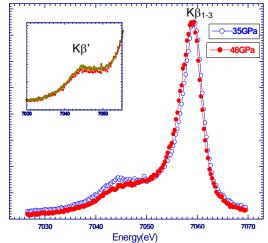
WehaveperformedXES measurements of the K  $\beta$  emission in Feup to pressures of 50 GPa. In the low pressure phase (ferromagnetic), in addition to the main, narrow K  $\beta^{1-3}$  peak, we observe a low energy satellite (K  $\beta$ ) – associated with transitions into the  $\beta$  or  $\beta$  -hole is a ligned opposite to the magnetic moment of the valence band).

At pressures between 12 -25 GPa, we observe the gradual disappearance of the K  $\beta$ 's satellite, indicating an electronic transition to a low -spinsta te. Our data shows that the quenching of the magnetism in Feocurs concurrently with the structural transformation to the high pressure (hcp) structure (as determined from

diffraction experiments). This result agrees with the reportbyRueffetal.[3].

Inourfirstexperiment, we observed a small but measurable positive shift in energy of the main K  $\beta$  peak with applied pressure (~1eV/GPa). However, our subsequent careful measurements, using remotely controlled membrane -type DACs to avoid sample motion in the beam, did not corroborate the original results. Comparison of the two experiments suggests that the error in the original result was caused by minor—but systematic—changes in beam position on the sample during pressure increase. Our analysis also shows that beam position changes do not affect the intensity of the  $K\beta$ ' satellite band, and therefo—re, do not affect the conclusions relating to the spin state of the valence band.

Figure 1 shows the results of our XES measurements under pressure on magnetite (Fe  $_3O_4$ ) using membrane DACs with Be gaskets. Our results show that, between 42GPa and 48GPa, the intensity of the K  $_\beta$ ' satellite (a measure of valence band spin polarization) decreases, stabilizing at a non-zero value. We interpret this to



**Fig.1:** XES measurements of the K bemission in magnetite (Fe <sub>3</sub>O<sub>4</sub>) at high pressure. Starting around 38GPa, the intensit y of the K β band decreases, indicating an electronic transition.

indicateanelectronic transition to an "intermediate spin" state. Further increases in pressure did not reduce the residual intensity of the Karaman state of the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the residual intensity of the Karaman state. Further increases in pressure did not reduce the reduce the reduced by th

This results tands in contrast with the published XES data of hematite (Fe  $_2O_3$ ) under pressure (and with our own measurements of Fe  $_2O_3$ ): in hematite, a strong suppression of the K  $\beta$ ' line in magnetite suggests that the final electronic state is a "low spin" state, with paired electronspins in the valence band [8].

### **Discussion:**

Our analysis of the observed changes in K  $\beta$  line -shape in Fe  $_3O_4$  under pressure suggests the dissociation of magnetite into FeO and Fe  $_2O_3$ . The result in ghigh -pressure phase of hematite is in a low spinstate, while the FeO maintains the polarization of the valence band at least to 143 GPa[9]. The result is an "intermediate -spin" state of the dissociated magnetite, with the observed reduced -intensity K  $\beta$ ' emission originating from the divalentiron sites. Above 50 GPa, the contribution of the low spin state of Fe to the total spin, is negligible.

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### **References:**

- 1.P.Soderlindetal., Phys.Rev.B50,5918(1994); Phys.Rev.B53,14063(1996).
- 2.H.K.Mao,etal.J.Appl.Phys. **38**,272(1967)
- 3. J. P. Rueff, M. Krisch, Y. Q. Cai, A. Kaprolat, M. Han Verbeni, and F. Sette, Phys. Rev. B **60**,14510 -14512(1999)
- 4.C -S. Yooetal, Phys. Rev. Lett, **84**,4132 -4135(2000)
- 5.M.Cohenetal, Science **275**,657(1997)
- 6.X.Wang,F.M.F.deGroot,andS.P.Cramer,Phys.Re v.B **56**,4553(1997).
- 7.R.J.Hemley,H.K.Mao,G.Shen,J.Badro,P.Gillet,M.Hanand,D.Haüsermann,Science (1997).
- 8. J. Badro, G. Fiquet, V.V. Struzhkin, M.S. Somayazulu, H.K. Mao, G. Shen, and T. Le Bihan. Phys.Rev.Lett. **89**,205504(2 002).
- 9. J. Badro, V.V. Struzhkin, J. Shu, R.J. Hemley, H.K. Mao, C. -C. Kao, J. -P. Rueff, and G. Shen. Phys.Rev.Lett. **83**,4101(1999).